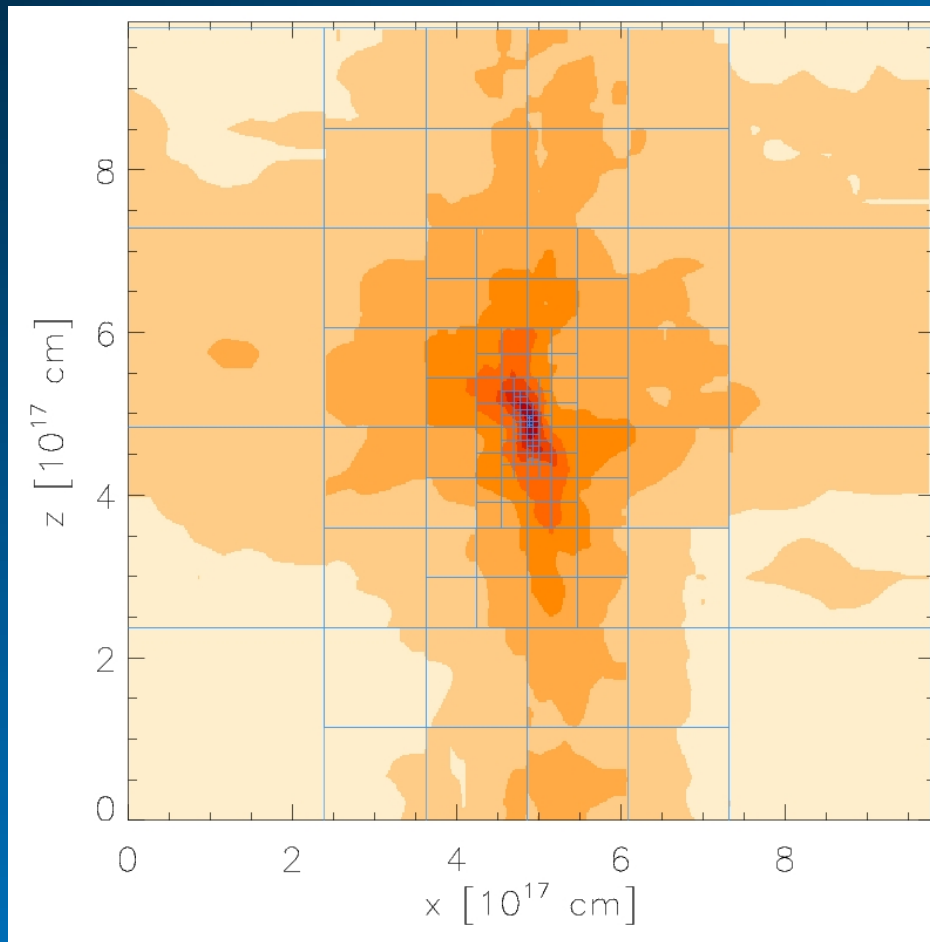


Modelling Star Formation with AMR*

Simulations

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ITA



3D MHD, AMR
Simulations of
collapsing cloud cores

*Adaptive Mesh Refinement

ITA Colloquium, Nov. 6, 2006

Outlook

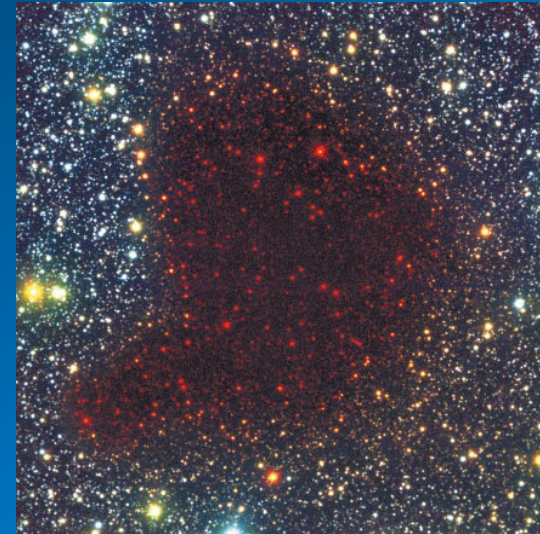
- Motivation / Introduction
- Numerical Method, AMR
- Modelling Cloud Cores
- Hydro Collapse Simulations
- ... with Magnetic Fields
- Outflows and Jets
- Turbulence
- Massive Stars
- Ongoing and Future Projects

Introduction

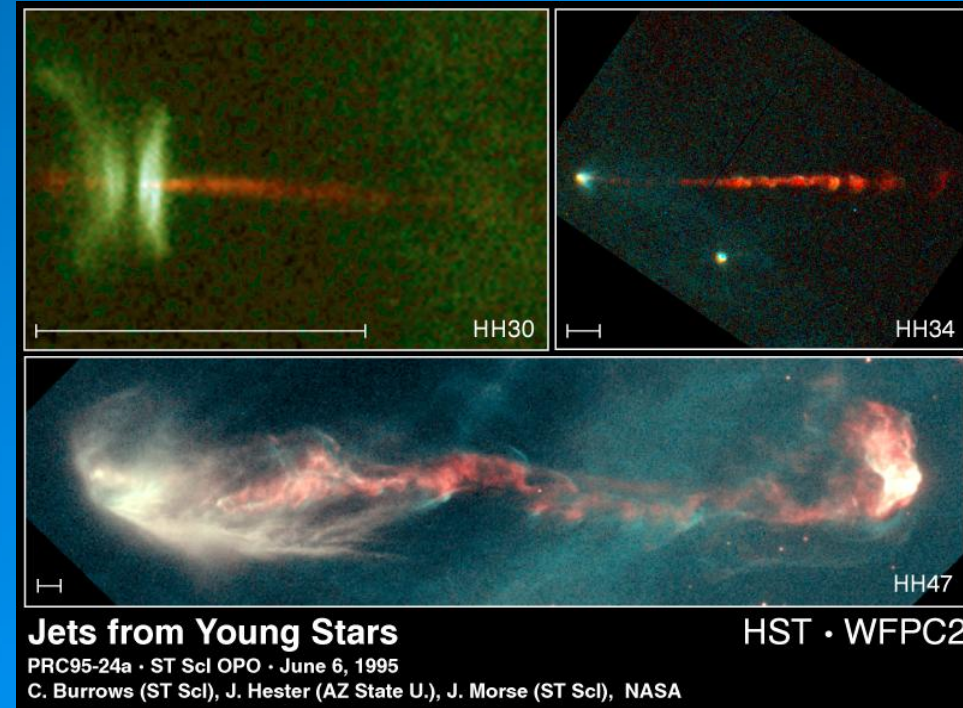
- Present day Star Formation in (Giant) Molecular Clouds (GMCs)
- Stars form out of collapsing cloud cores



Orion Nebula (M 42), Star Forming region (*HST image*)



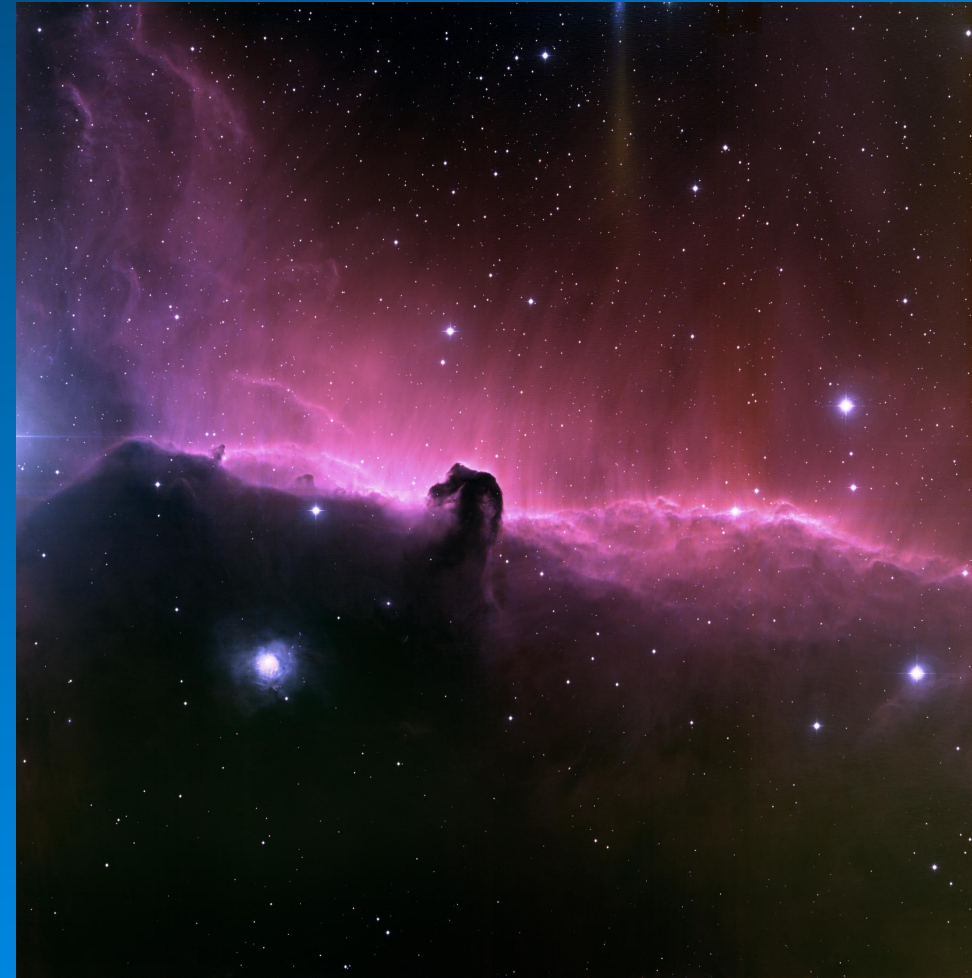
Barnard 68,
Cloud core (cold,
self shielded)
(*Alves, Lada & Lada,
Nature 2001*)



Motivation for numerical simulations

Problems

- Self **Gravity**: non-linear interactions
- Non-vanishing **angular momenta** \Rightarrow disk formation
- Chemical evolution, **Cooling** & Heating
- **Magnetic Fields**: Jets and Outflows?
- Radiation Feedback

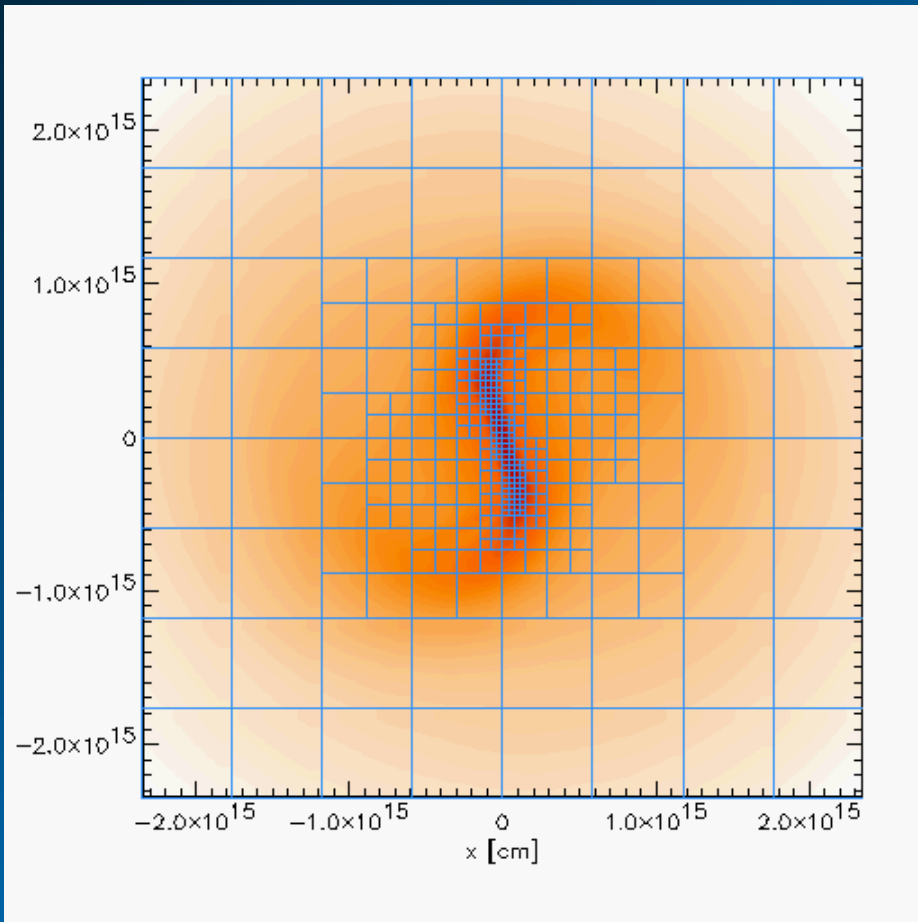


Horsehead Nebula (Barnard 33) in the OMC
NOAO image

Solution \Rightarrow **Direct Numerical Simulations**

Numerical Method

FLASH* Code



- 3D **Grid** based MHD integrator for parallel computing (MPI)
- Hydro solver: PPM, Kurganov
- MHD solver: MUSCL (van Leer) with $\text{div}\mathbf{B}$ cleaning
- Gravity: multigrid or multipole, periodic or isolated BCs
- **AMR**: block structured (PARAMESH library);
block resolutions vary by factors of **two**
- Refinement on own choice (e.g. gradient, curvature, density, etc.)
- IDL routines for visualization

Pros

- modular, easy to use
- large community: e.g. multi fluid nuclear reactions, RT module, N-body particles, cosmology
- support from developers

Cons

- resource consuming
- not very fast
- block structured AMR (will be improved with newer versions)

**Alliance Center for Astrophysical
Thermonuclear Flashes (ASC),
University of Chicago
Current Version: 3.0*

Numerical Method

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) &= 0 \\ \frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B}) + \nabla p_* &= \rho \mathbf{g} + \nabla \cdot \boldsymbol{\tau} \\ \frac{\partial \rho E}{\partial t} + \nabla \cdot (\mathbf{v}(\rho E + p_*) - \mathbf{B}(\mathbf{v} \cdot \mathbf{B})) &= \rho \mathbf{g} \cdot \mathbf{v} + \nabla \cdot (\mathbf{v} \cdot \boldsymbol{\tau} + \boldsymbol{\sigma} \nabla T) + \nabla \cdot (\mathbf{B} \times (\eta \nabla \times \mathbf{B})) \\ \frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v}) &= -\nabla \times (\eta \nabla \times \mathbf{B}) \end{aligned}$$

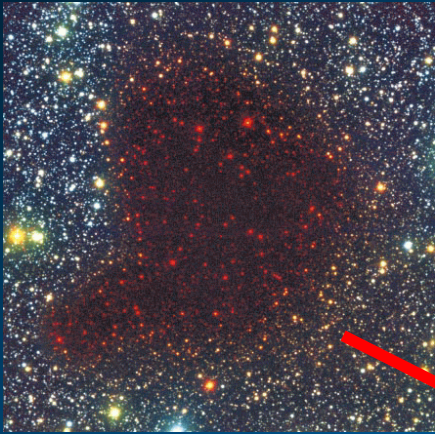
$$p_* = p + \frac{B^2}{2},$$

$$E = \frac{1}{2} v^2 + \varepsilon + \frac{1}{2} \frac{B^2}{\rho},$$

$$\boldsymbol{\tau} = \mu \left((\nabla \mathbf{v}) + (\nabla \mathbf{v})^T - \frac{2}{3} (\nabla \cdot \mathbf{v}) \right)$$

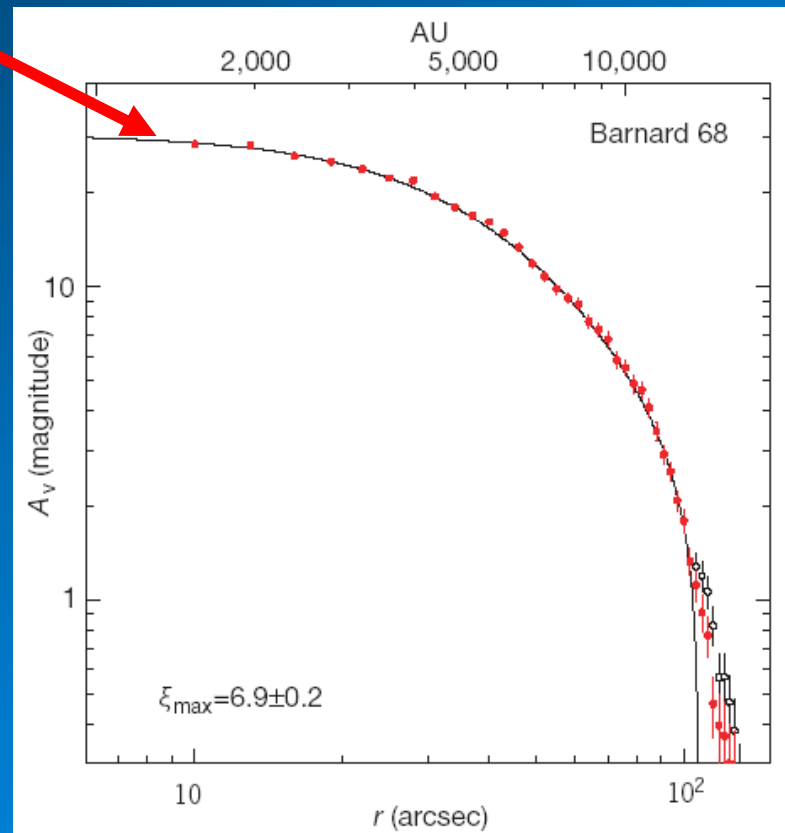
Applications

Collapse of Hydrostatic Cores



Bok Globule B 68

Dust column density profile in terms of visual extinction follows a BE-Profile
mass $\sim 2.1 M_{\text{sol}}$



Molecular Clouds in hydrostatic equilibrium follow a **Bonnor-Ebert-Profile**

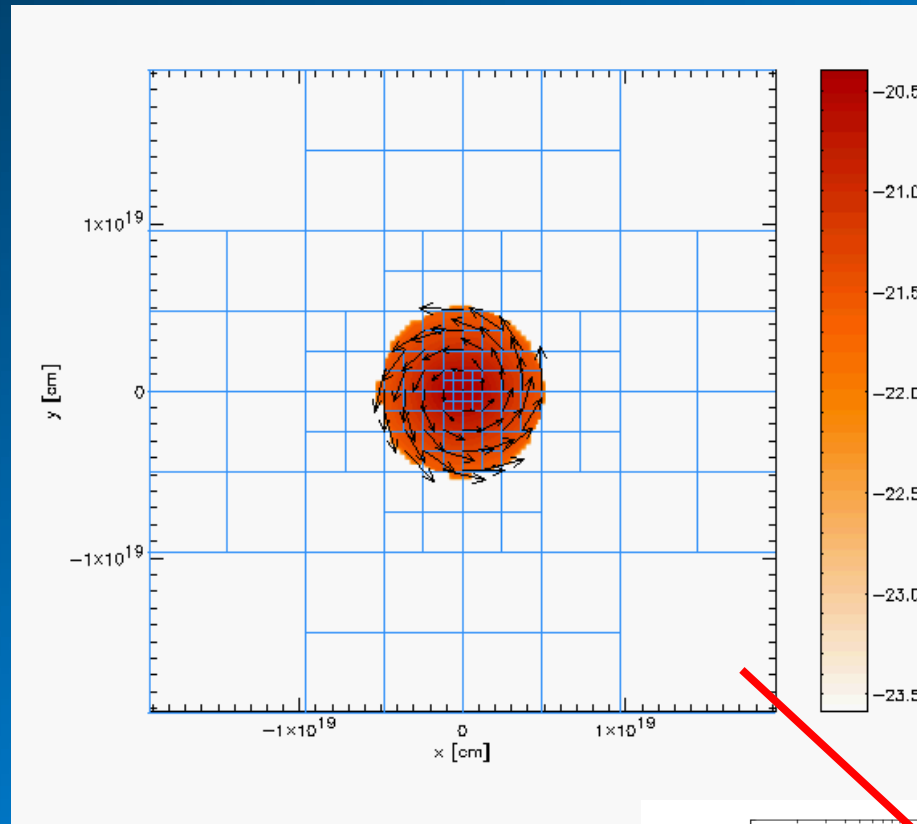
Critical BE Sphere:
 $\xi = 6.451$

Other observational evidences:

- Coalsack globule 2,
 $M \sim 4.5 M_{\text{sol}}$, (*Racca, Gomez & Kenyon, 2002*)
- Dark globule B335,
 $M \sim 14 M_{\text{sol}}$,
(*Harvey et al. 2001*)

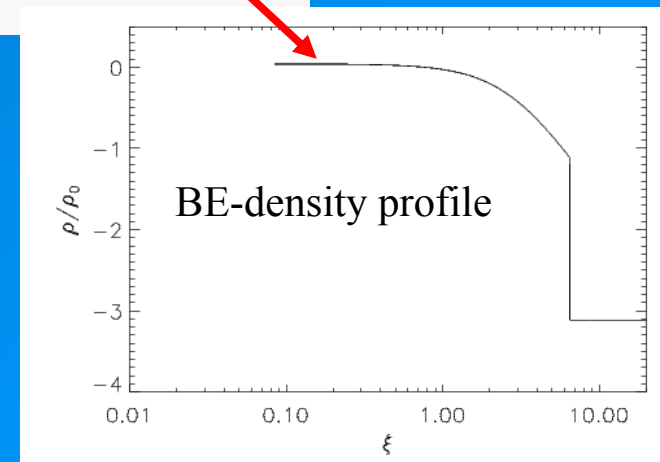
Collapse of Hydrostatic Cores

- Slowly rotating Bonnor-Ebert-Spheres
- Low Mass $M \sim 2.1 M_{\text{sol}}$
- High Mass $\sim 170 M_{\text{sol}}$
- **Cooling** due to molecular excitations, gas-dust interaction, H_2 dissociation
- AMR \Rightarrow resolves **Jeans length** with more than 8 grid points during collapse (*Truelove et al. 1997*)
- Up to 27 refinement levels (dynamical range $\sim 10^7$)

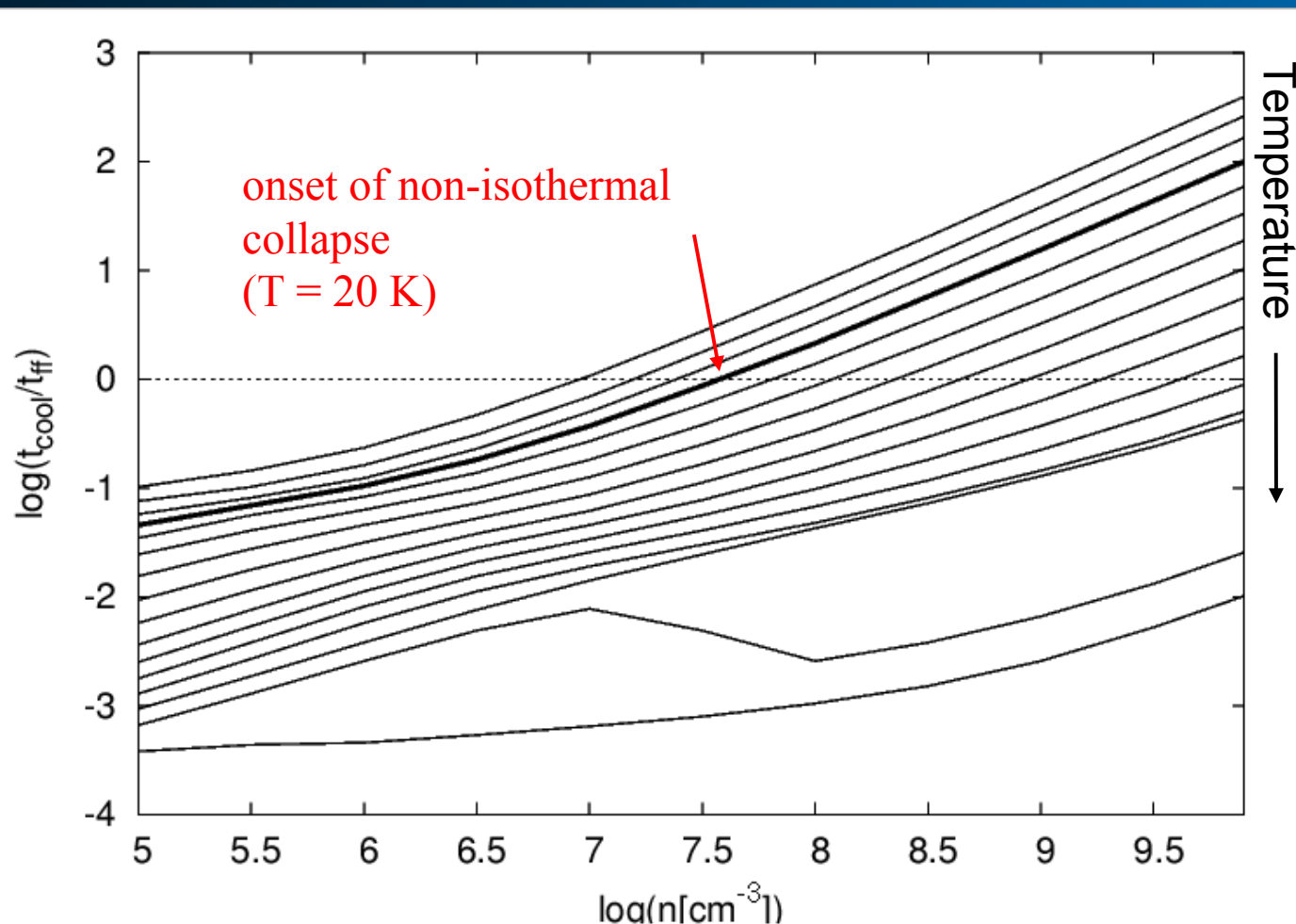


Initial conditions:

- cool molecular cloud ($T = 16 \text{ K}$)
- hot ambient, low density, medium (pressure match at the sphere boundary)
- $\Omega t_{\text{ff}} = 0.1 - 0.4$



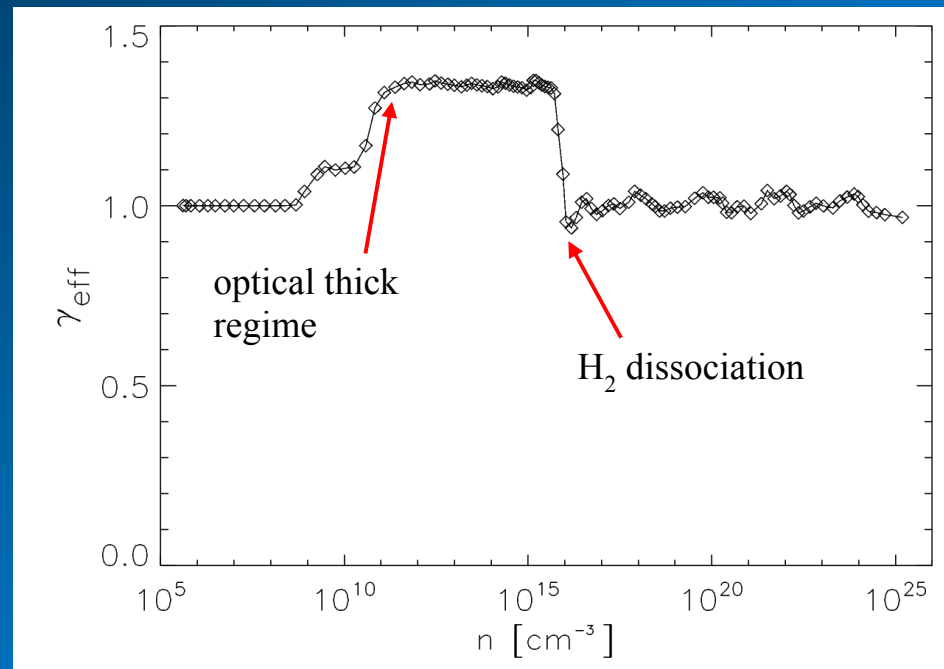
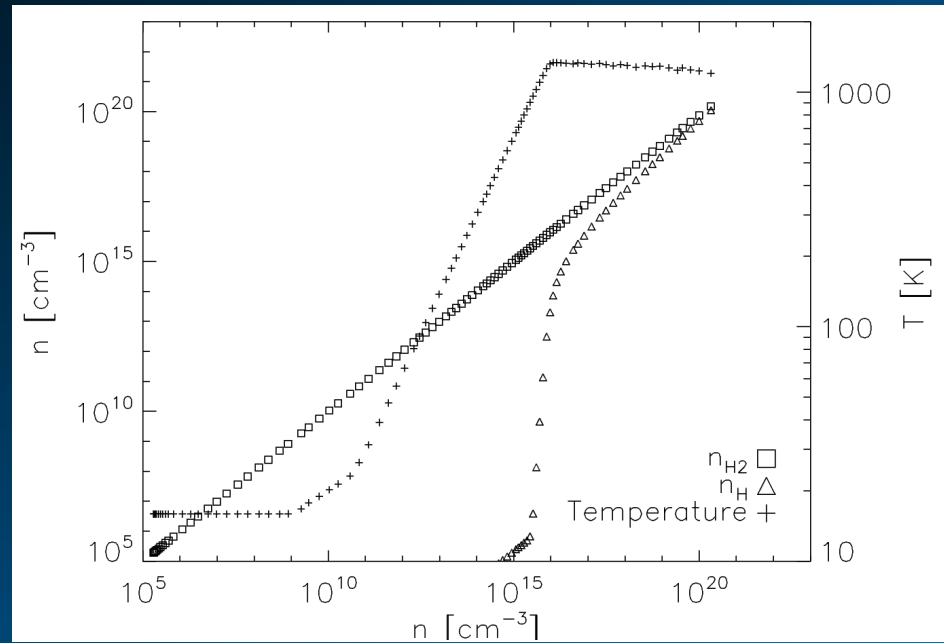
Effect of molecular line cooling



- Main coolants: H_2O , CO , H_2 , O_2
- Efficient cooling in low density regime $n \sim 10^{7.5}$ **(no dust)**
 \Rightarrow isothermal collapse
- Cooling sets a fixed physical length scale, $r \sim c/(G\rho_{\text{crit}})^{1/2}$ where $t_{\text{cool}} \sim t_{\text{ff}} \Rightarrow$ warm core and **shocked** gas

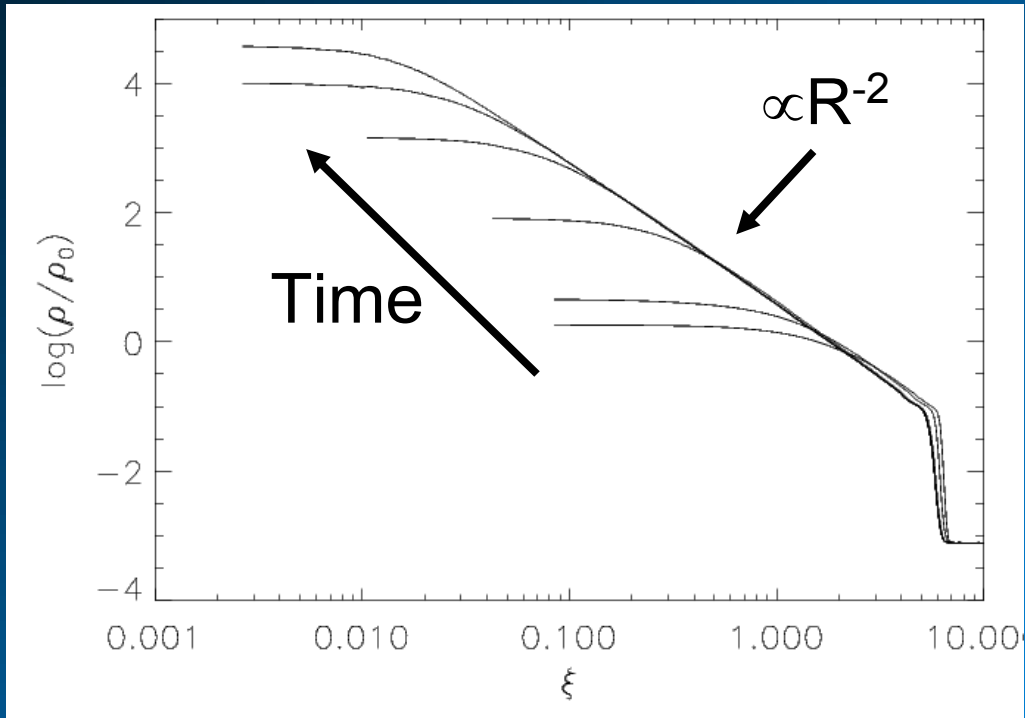
Cooling data from *Neufeld & Kaufman, ApJ, 1993, and Neufeld, Lepp & Melnick, ApJS, 1995*: radiative losses due to collisional excitations

Cooling with Dust & H₂ dissociation

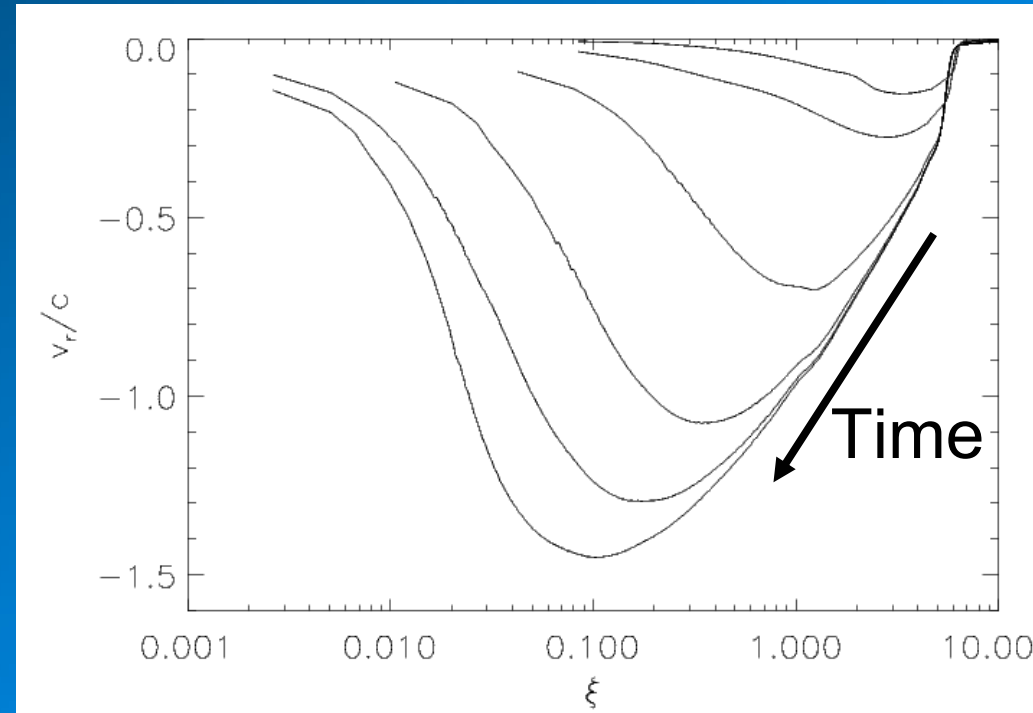


- **Dust**-gas interactions (*Goldsmith 2001*) keeps the gas isothermal until $n \sim 10^{11}$ cm⁻³ \Rightarrow scale of hot core: \sim few x 10 AU
- **Optically thick** at $n \sim 10^{11}$ cm⁻³ \Rightarrow heating with $T \sim n^{1/3}$
- **H₂ dissociation** at ~ 1200 K (*Shapiro & Kang 1987*) \Rightarrow isothermal collapse (second collapse; *Larson 1969*)
- dissociation process is “self-regulating” due to strong temperature dependence

Isothermal Bonnor-Ebert collapse



density



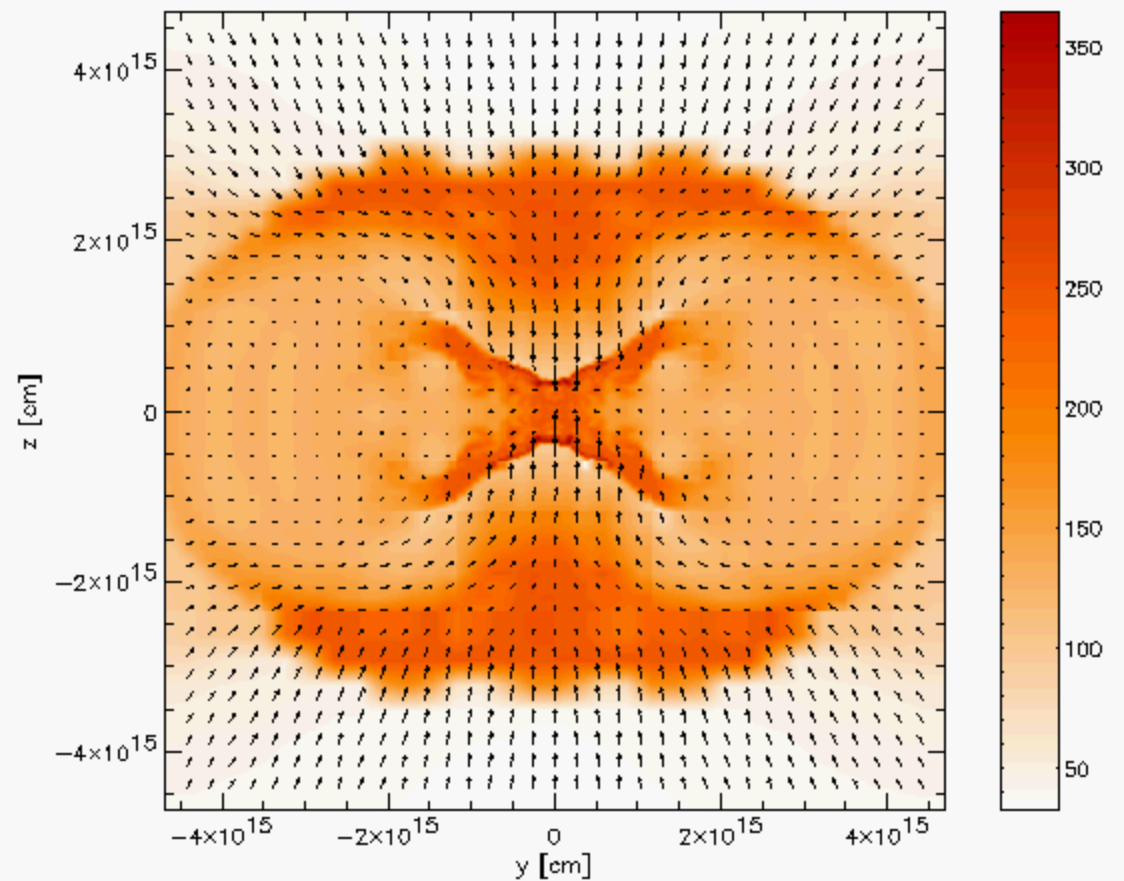
radial velocity

Outside-in
non-homologous collapse

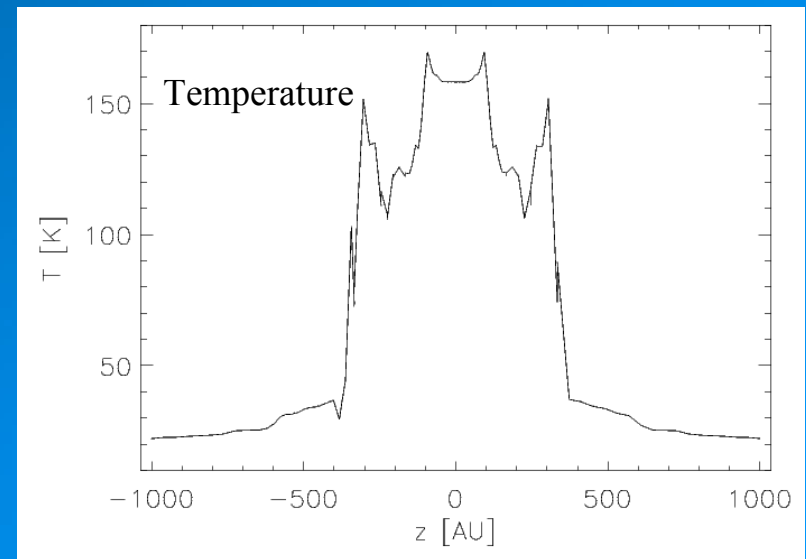
Non-isothermal collapse phase

Shock structures

Temperature



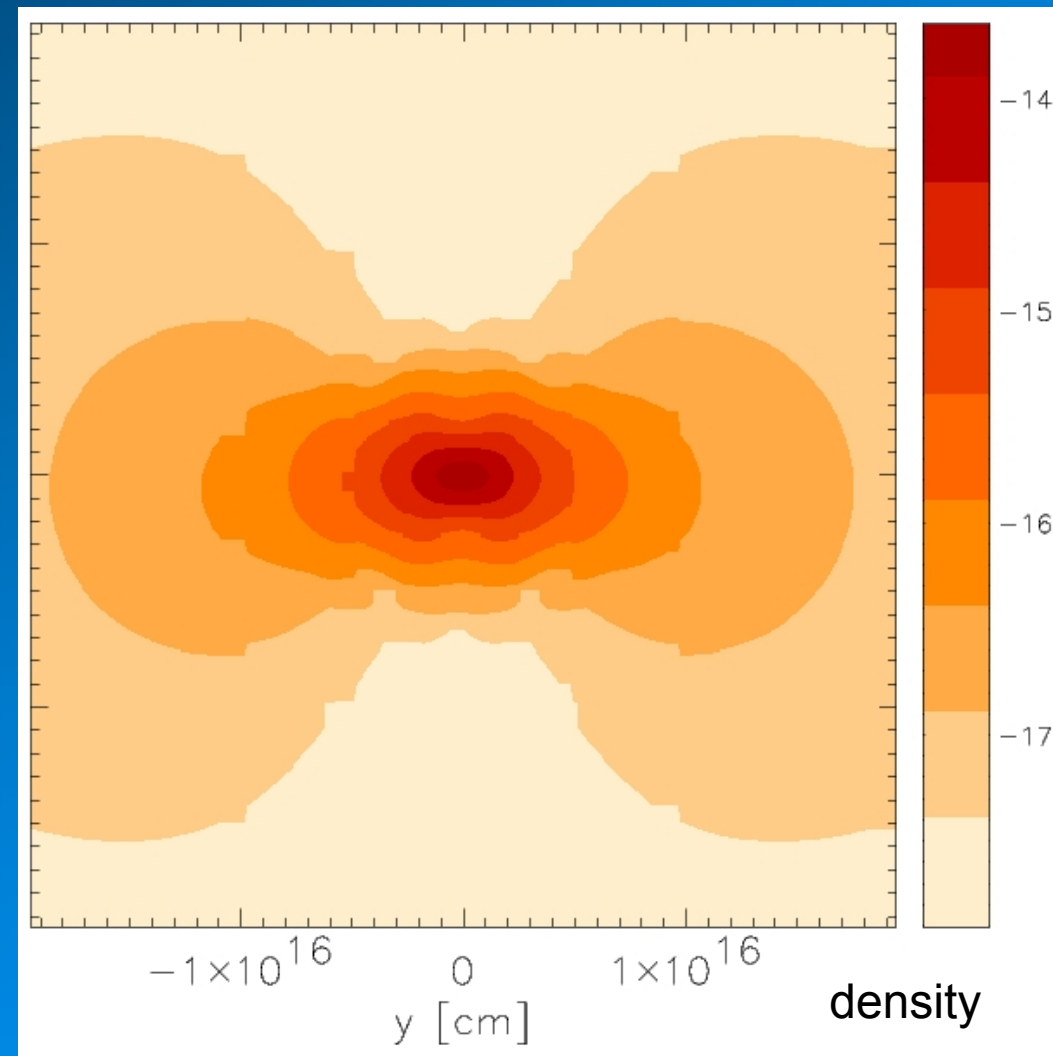
At least 2 shock fronts build up during the collapse with Mach numbers $\sim 2 - 4$



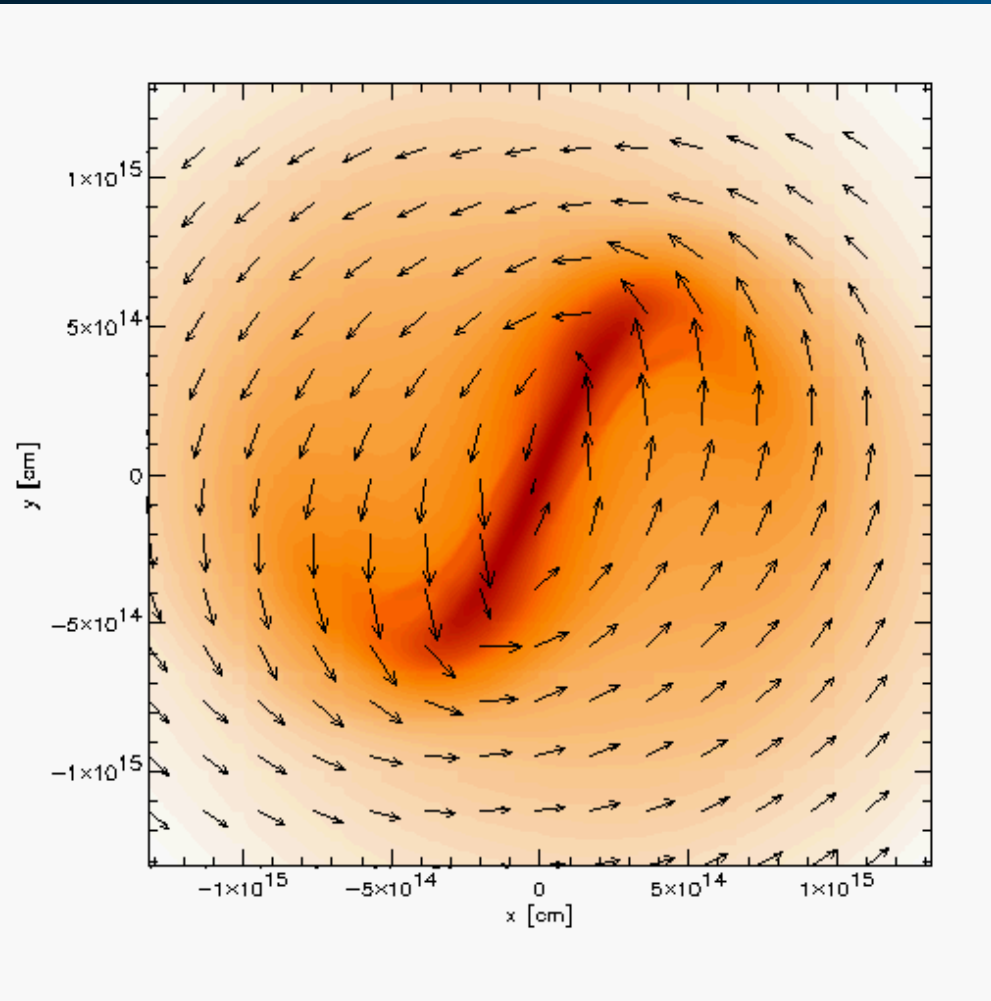
Note: only molecular line cooling
no dust-gas interaction

Disk Formation / Disk Structure

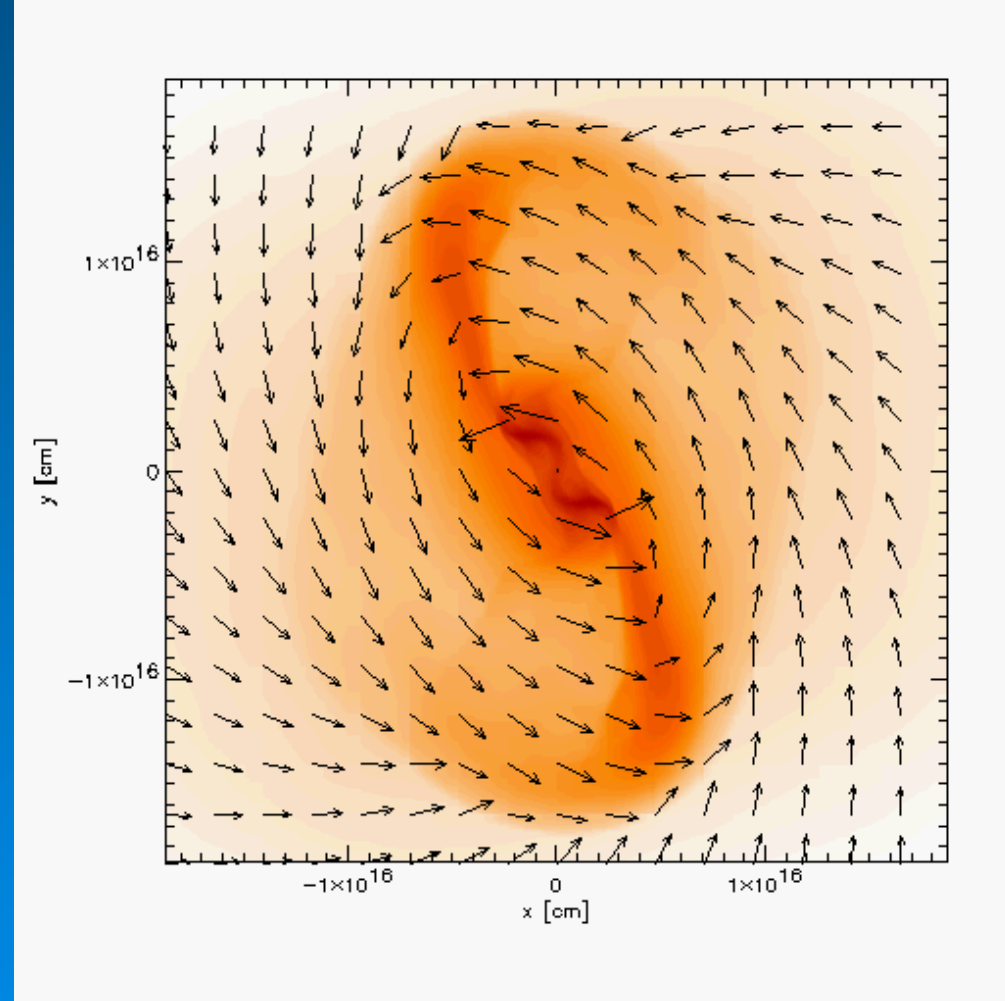
- Angular momentum conservation \Rightarrow formation of a **disk**
- Initial rotation determines disk size and structure (e.g. spiral arms, ring, binaries)
- Size of protodisk: few 100 AU – 10^3 AU depending on initial spin
- BE case: $\Omega f_{\text{ff}} = 0.1 \Rightarrow$ **bar** formation;
- $\Omega f_{\text{ff}} = 0.2 \Rightarrow$ **fragmentation**
(*Matsumoto & Hanawa 2003, Banerjee et al. 2004*)
- Transport (redistribution) of angular momentum due to spiral arms and magnetic fields



Disk Formation / Disk Structure

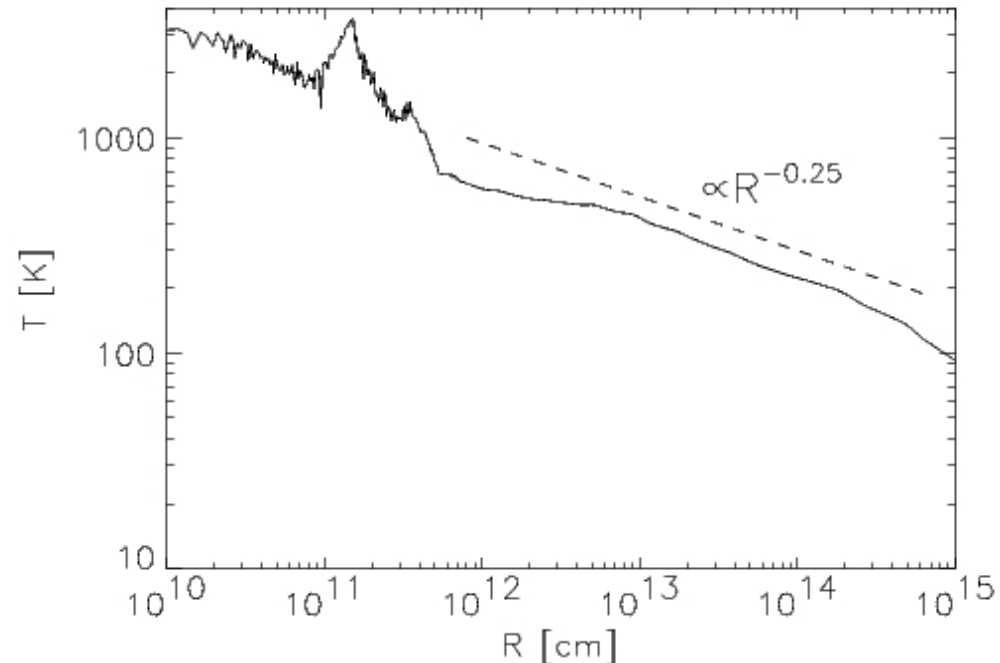
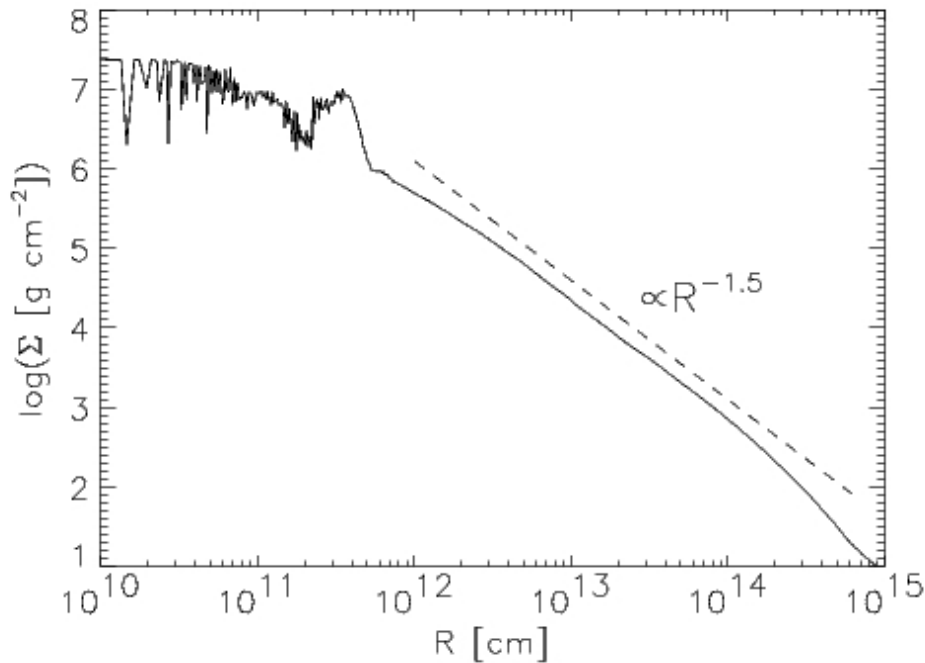


bar formation with
slow rotation



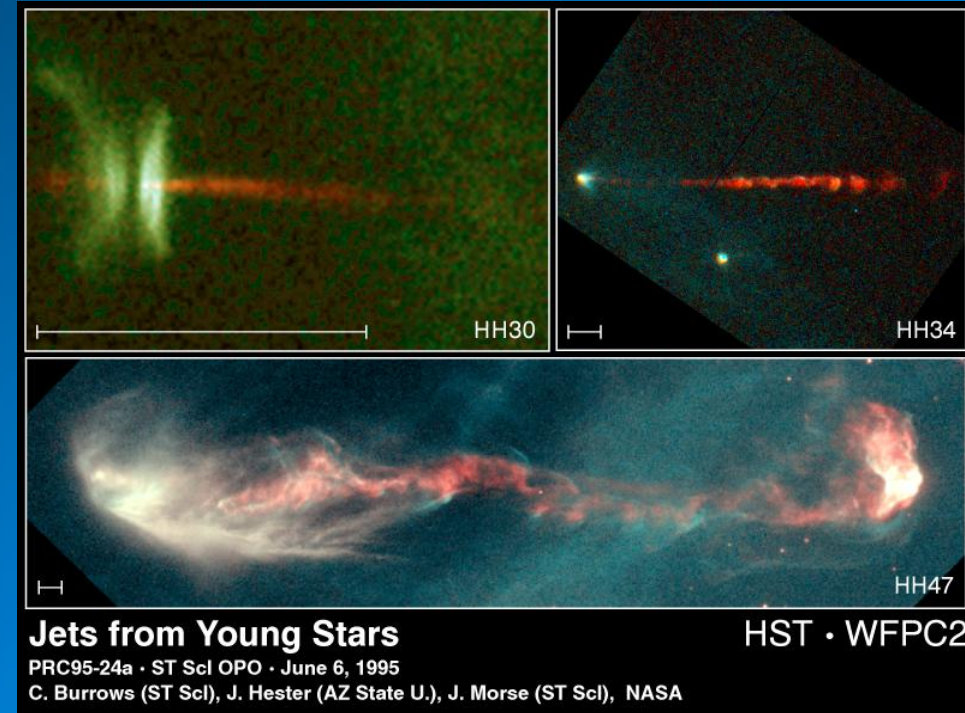
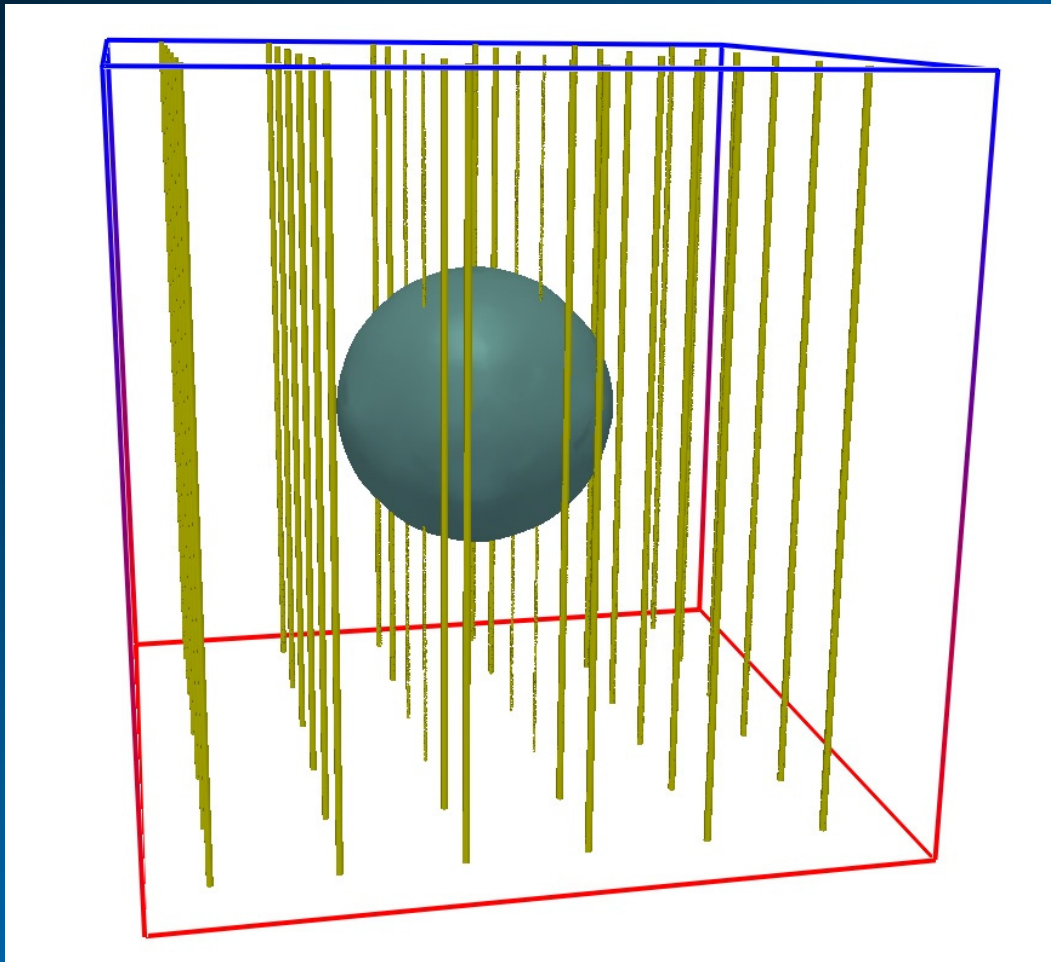
fragmentation with faster initial rotation

Disk structure



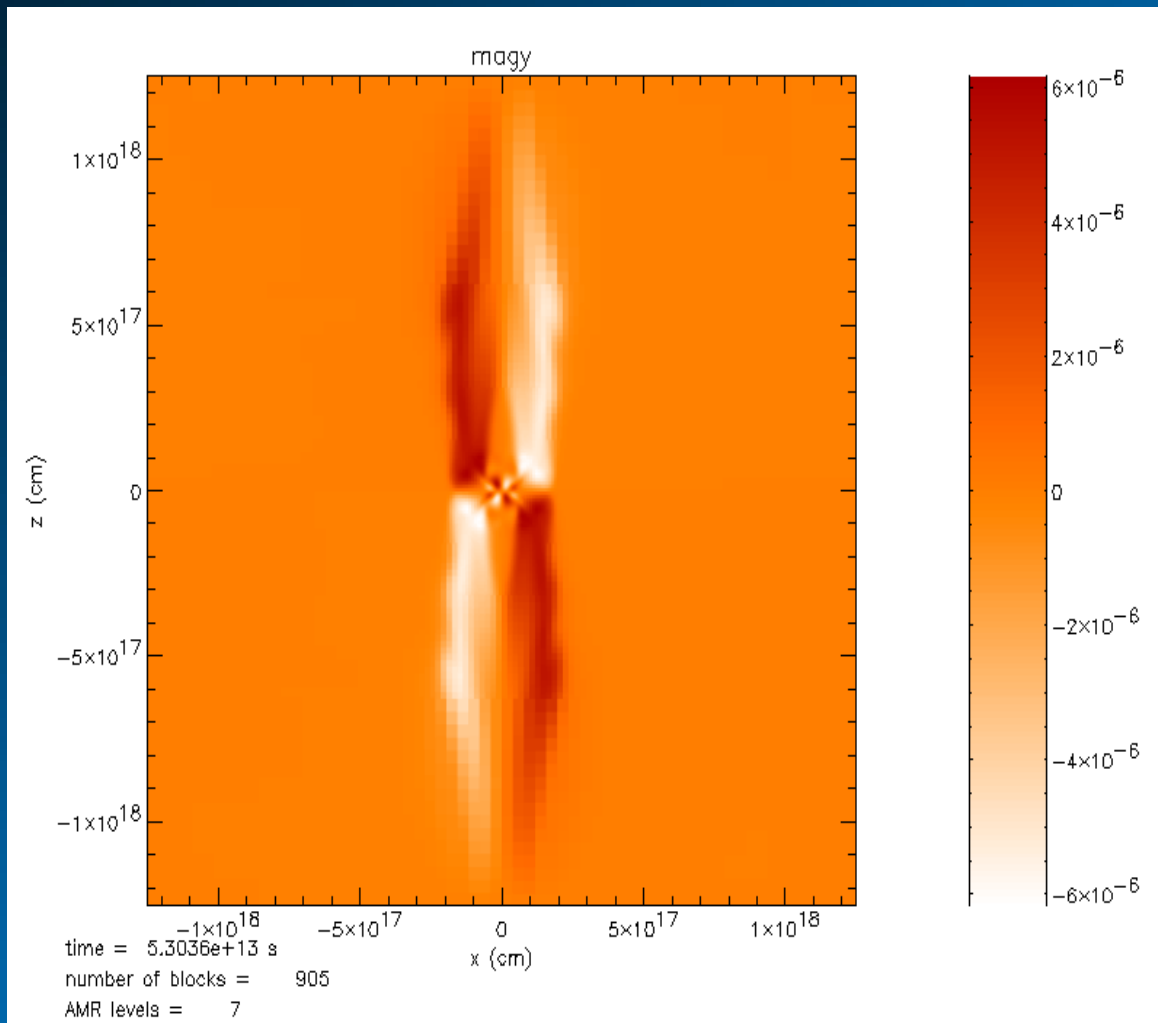
- Disk profile: Hayashi-type $\propto R^{-3/2}$ (Hayashi, 1981)
- Shallow temperature profile $\propto R^{-1/4}$

Magnetic Fields



- Jets / Outflow from YSOs magnetically driven?
- **Ideally** coupled to the gas (no ambipolar diffusion)
- Initially not dominant;
 $P_{\text{therm}}/P_{\text{mag}} \sim 80$; $B \sim \mu\text{Gauss}$

Pre-collapse, Magnetic braking

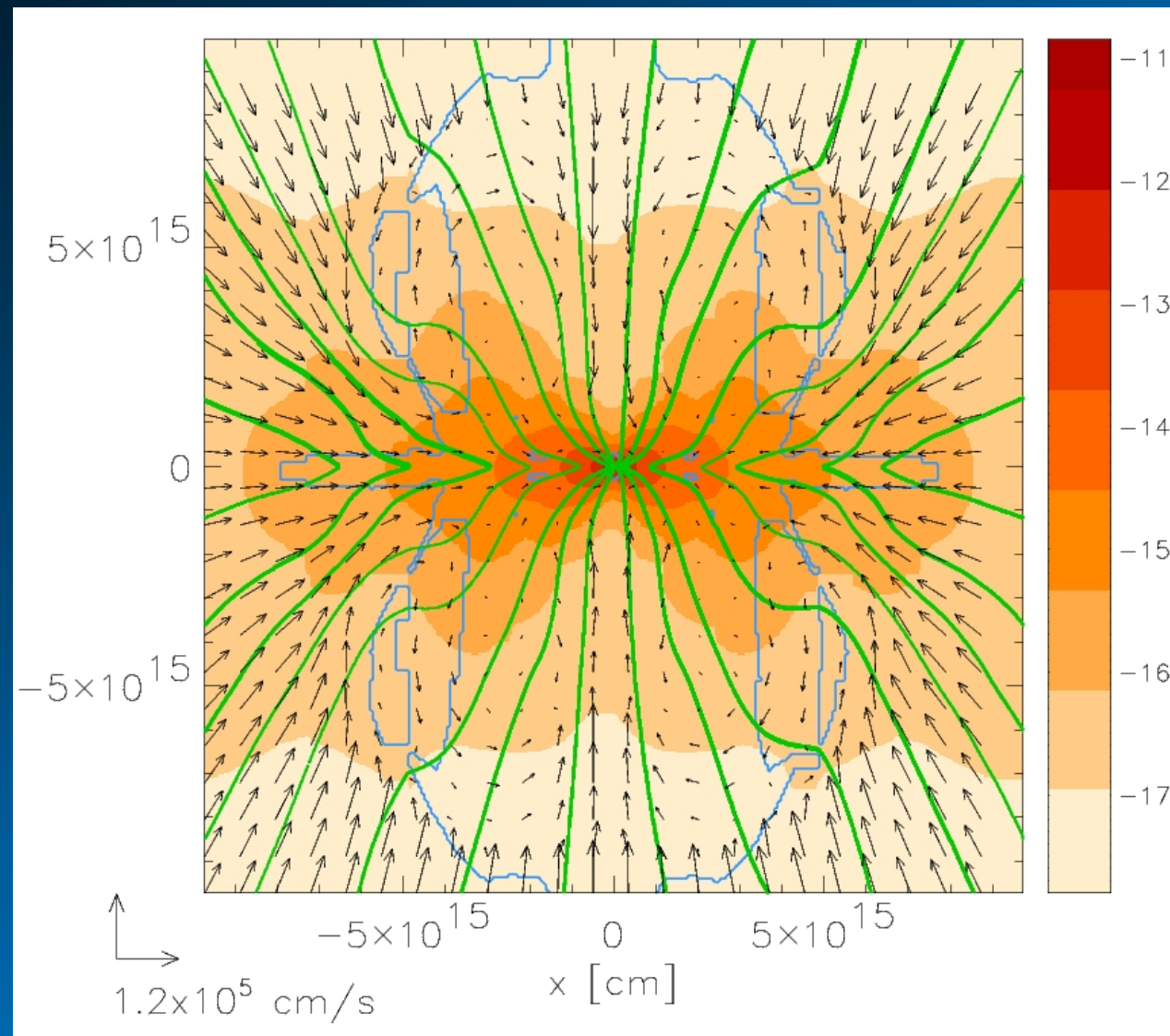


- Linear regime: torsional Alfvén wave is launched in the ambient medium (Poynting flux)
- Loss of initial angular momentum \Rightarrow **Magnetic braking** (*Mouschovias & Paleologou, 1980*)
- Spin down time:

$$\tau_{\text{damp}} \approx \frac{\pi R}{4 v_{A,\text{ext}}} \delta$$

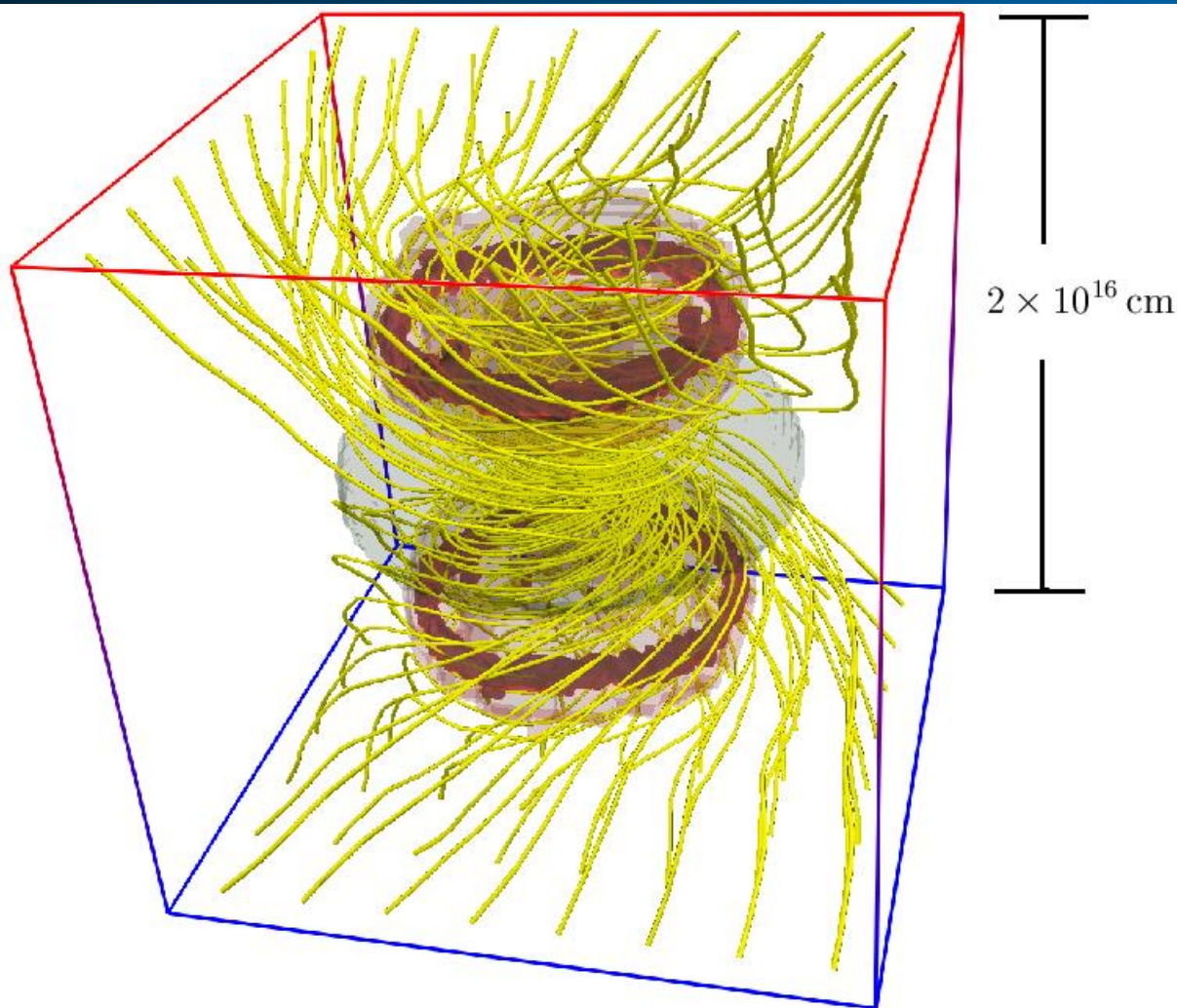
- In our case: $\sim 10^6$ years

Large scale outflow



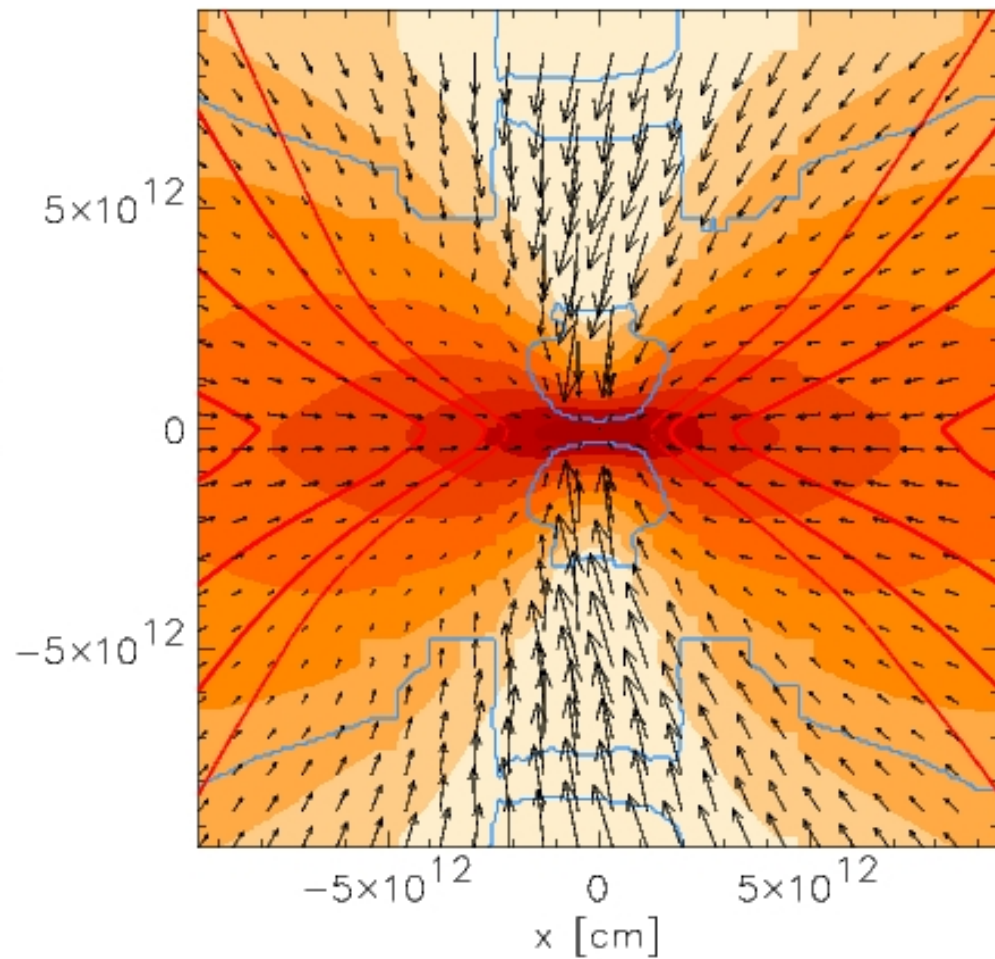
- Magnetic field is **compressed** with the gas
- Rotating disk generates **toroidal** magnetic field \Rightarrow **outflow** (*Blandford & Payne, 1982*)
- Shock fronts are pushed outwards (magnetic tower; *Lynden-Bell 2003*)
- Outflow velocities $v \sim 0.4$ km/sec, $M \sim 2-3$
- Accretion: funneled along the rotation axis, through disk

Large scale outflow

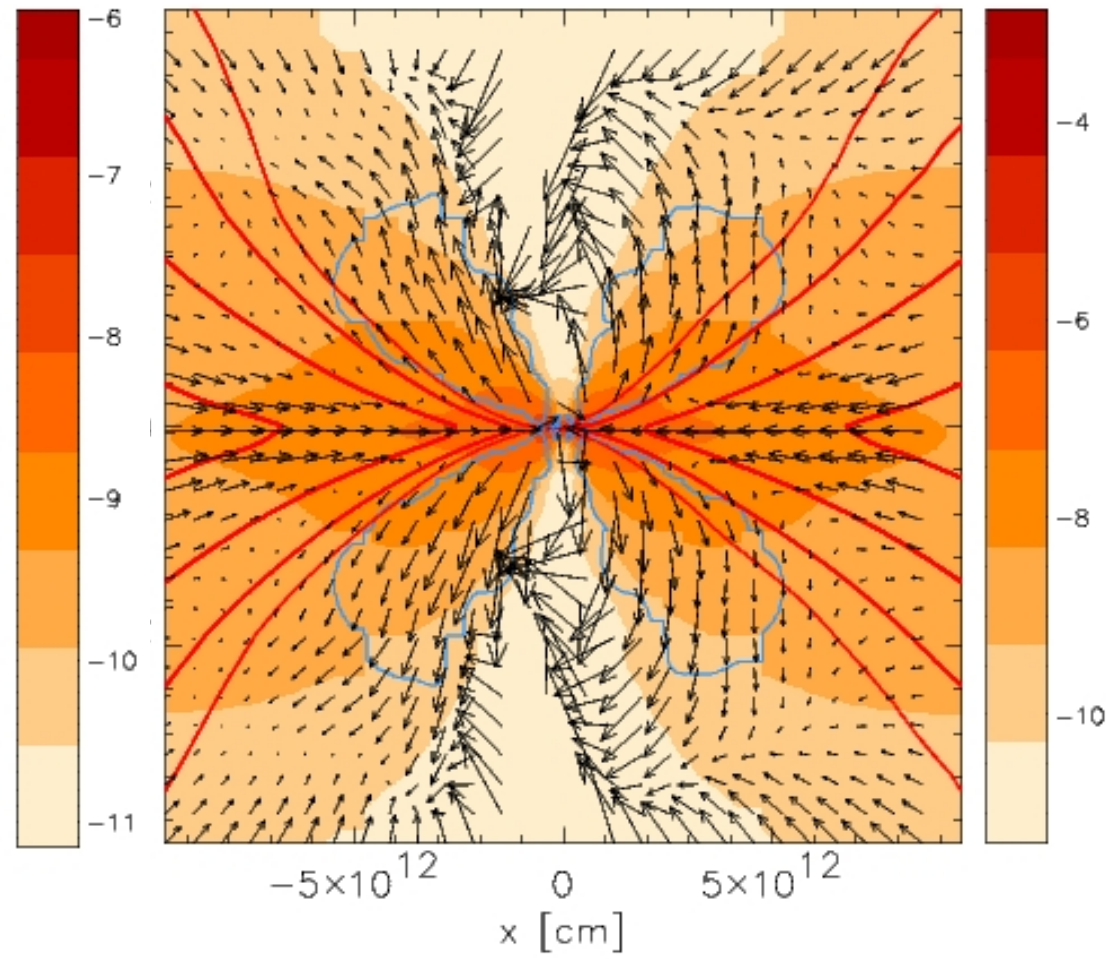


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Jets from the inner disk

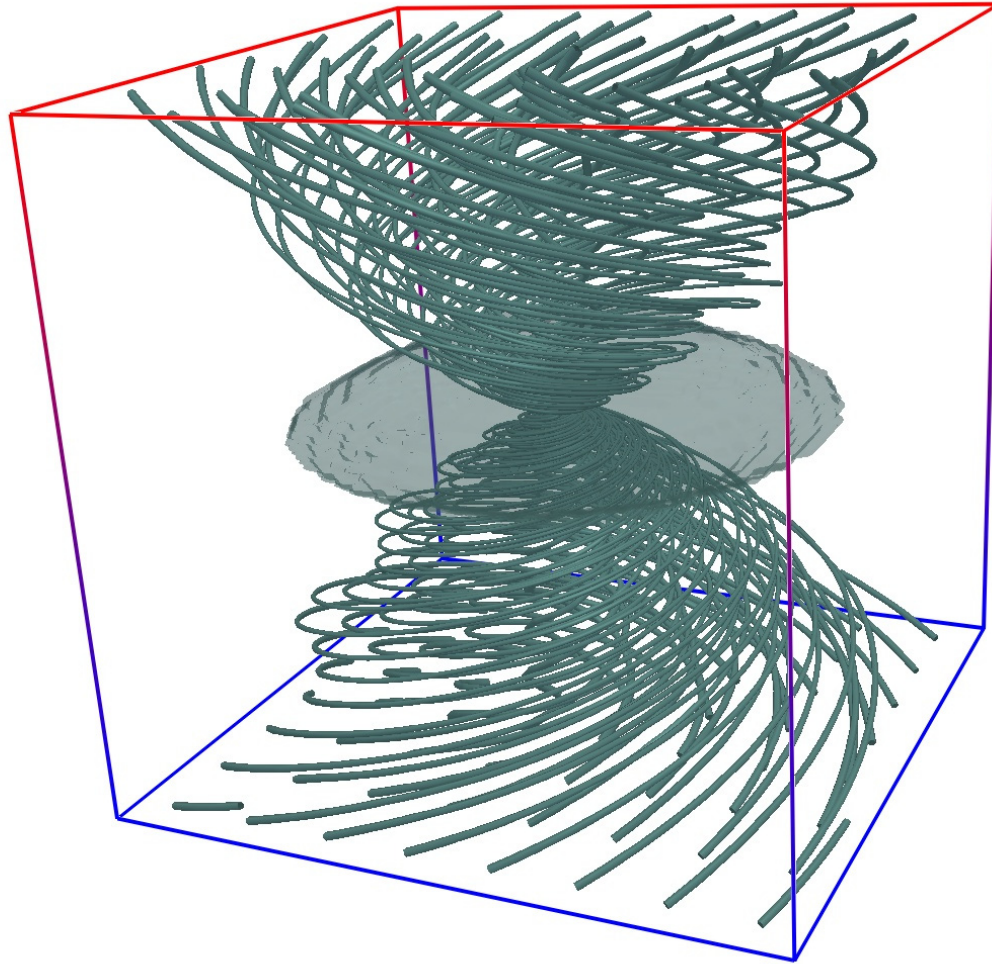


Inner disk, collapse phase



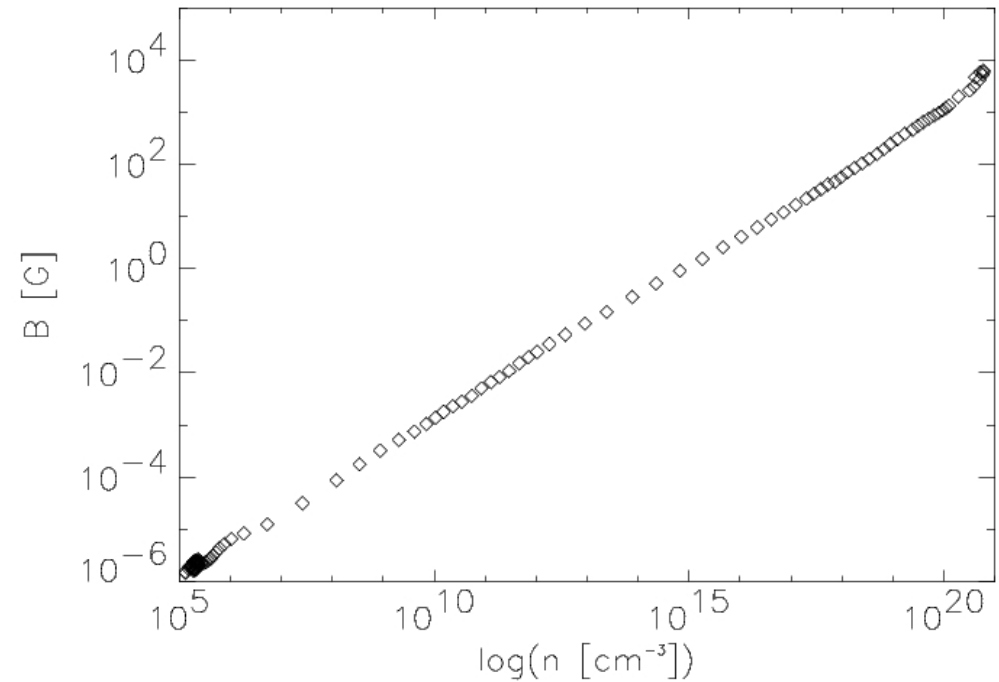
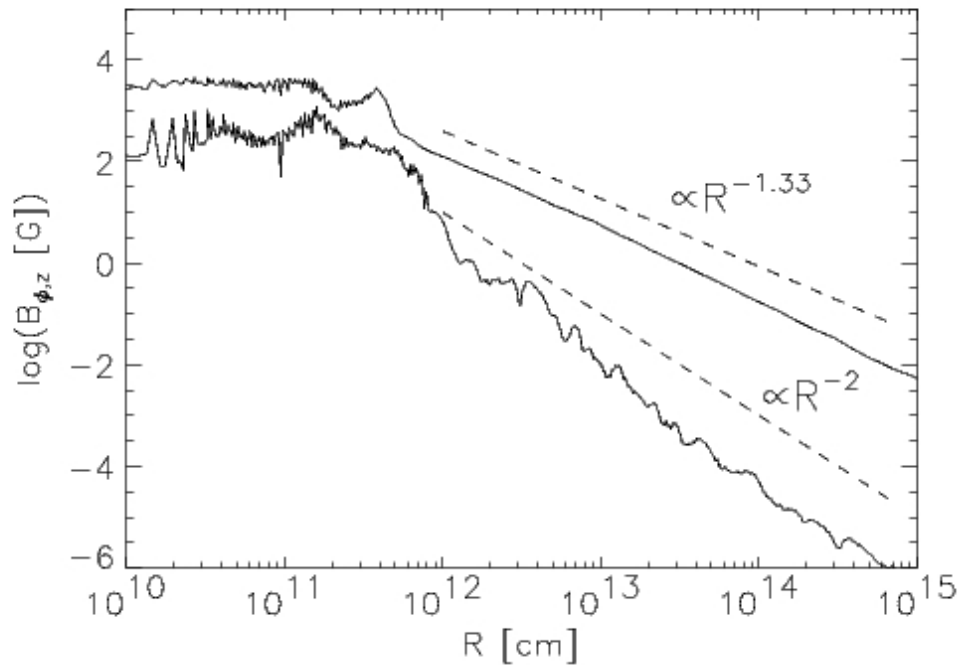
5 month later: a Jet is launched

Jets from the inner disk



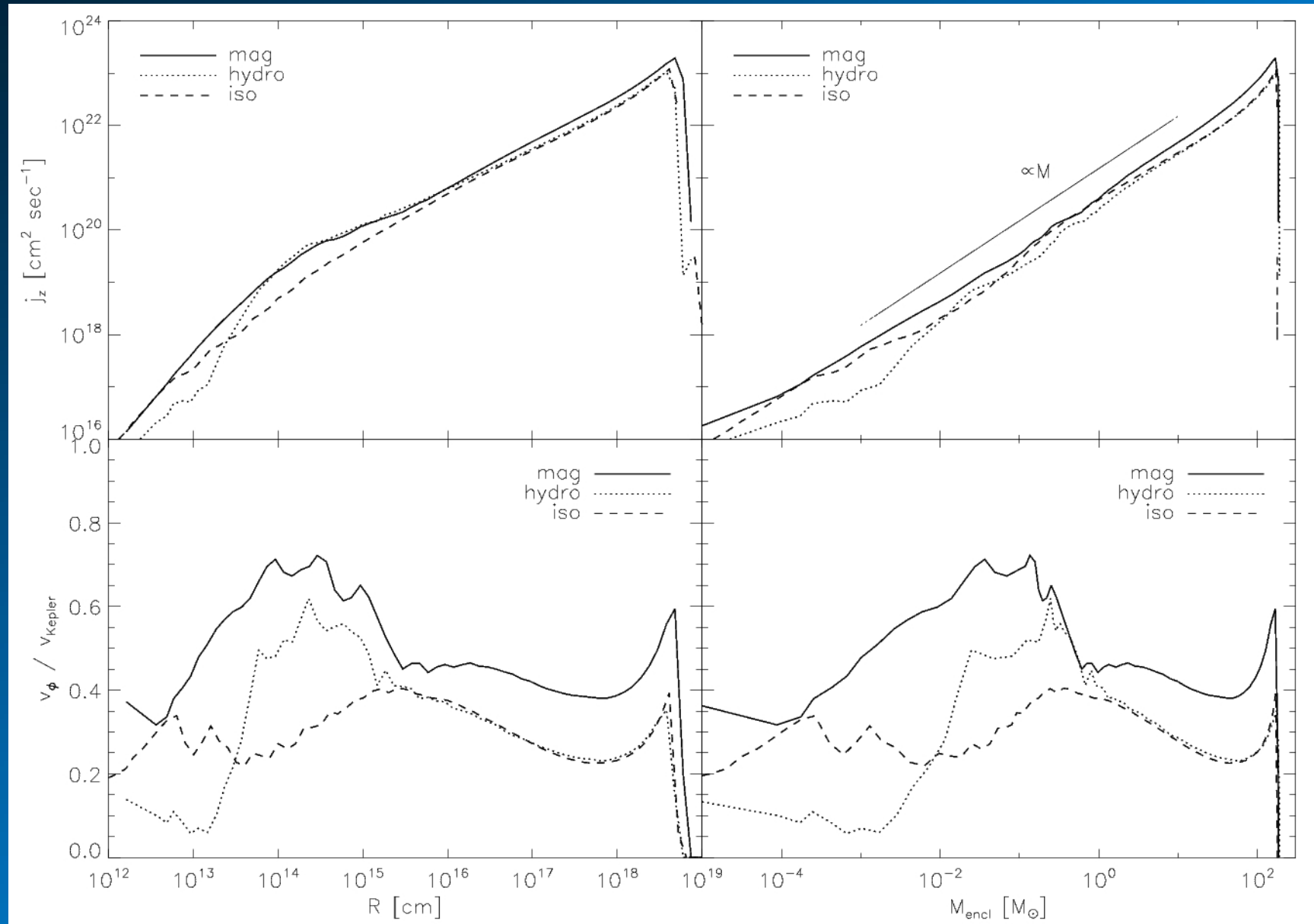
- Strongly pinched and warped field structure
- Jet velocities $v \sim 3$ km/sec
Super-sonic ($M \sim 4$),
Super-Alfvénic

Magnetic field structure / evolution



- $B_z > B_\phi$ in the core and disk (expectation from a stationary accretion disk $B \propto R^{-1.25}$; *Blandford & Payne 1982*)
- $B_{\text{core}} \propto n^{0.6}$
- Expected field strength in the protostar $\sim 10^4 - 10^5$ G
- Potential seed field for Ap stars (*Braithwaite & Spruit, 2004*)

Angular Momentum

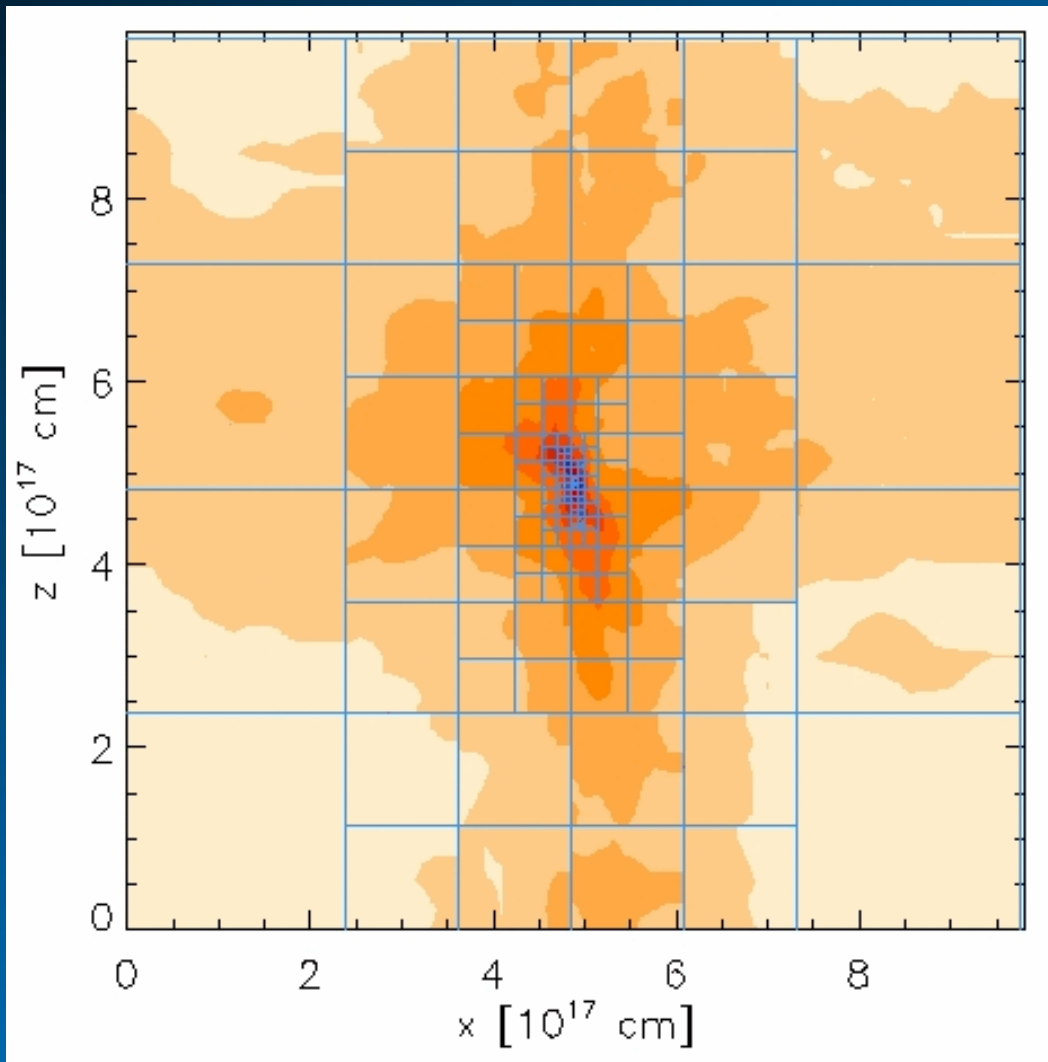


Turbulence



- Star Formation takes place in a **supersonic turbulent** environment (e.g. *Mac Low & Klessen 2004*)
- Hydrostatic cores??? (but *Shu et al. 1987, Mouschovias*)

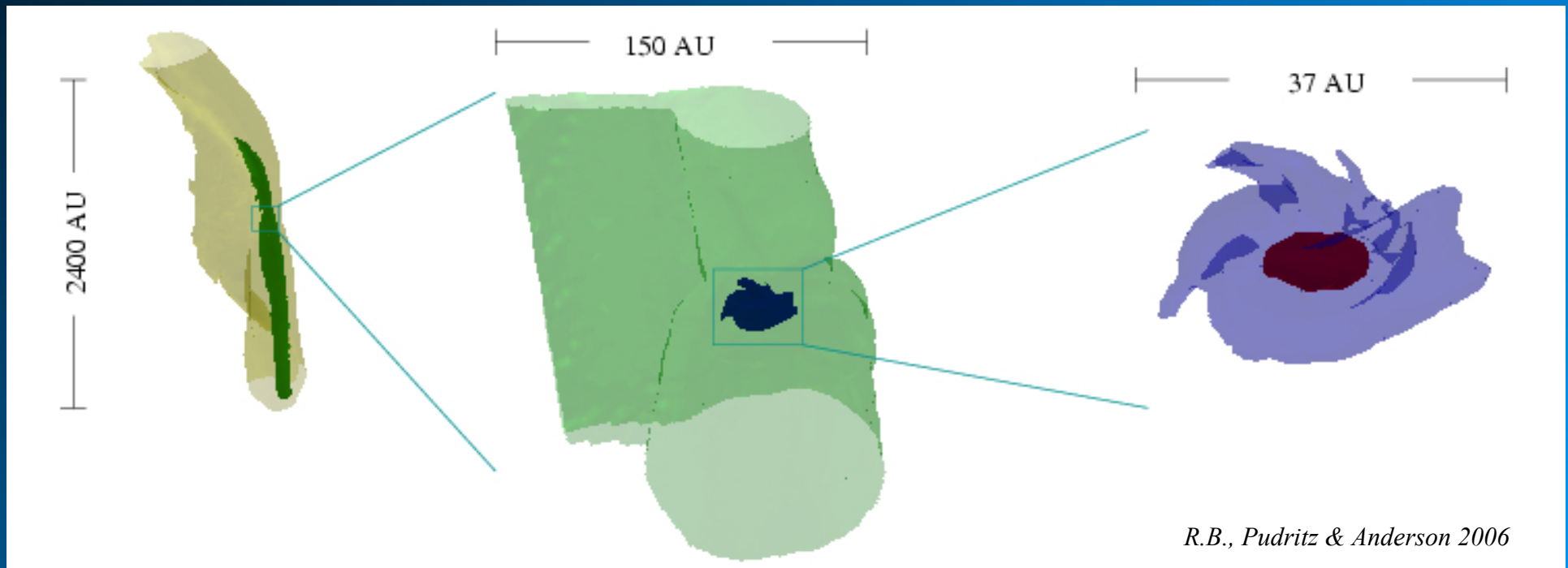
Collapse with supersonic turbulence



Initial setup as “seen” by the FLASH code

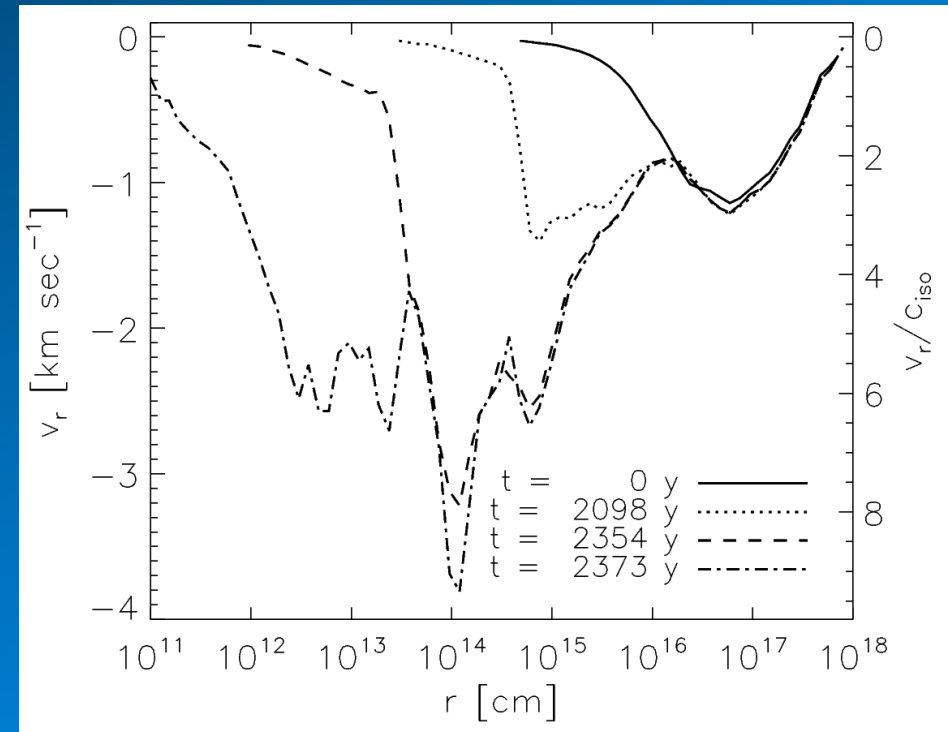
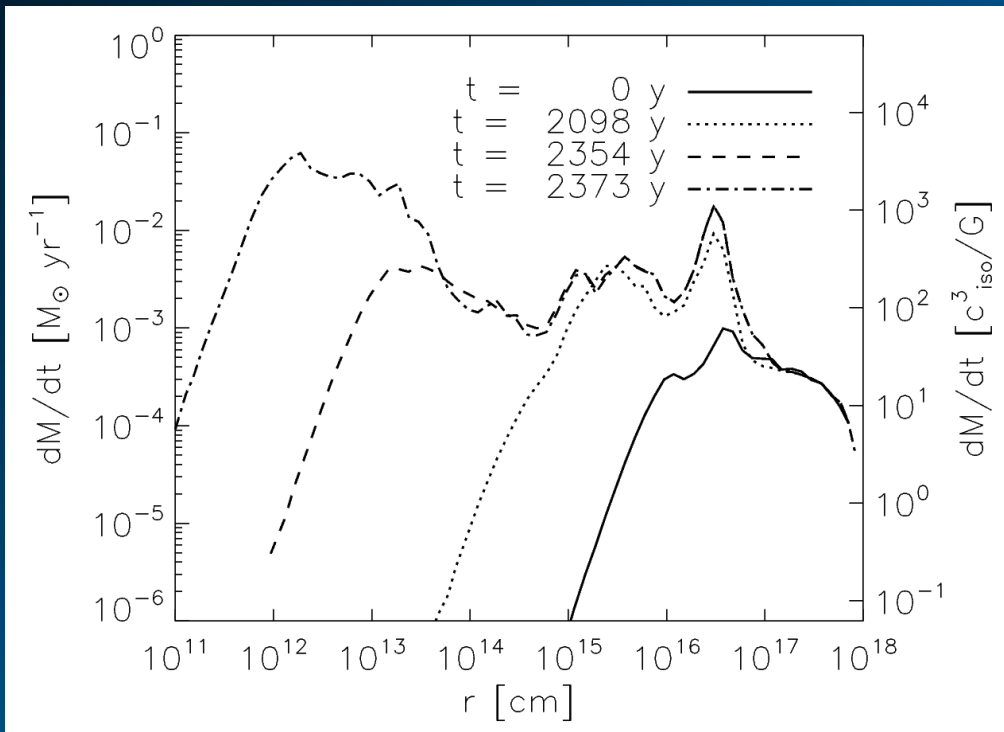
- Initial data from *Tilley & Pudritz 2004*: ZEUS simulations of core formation within a supersonic **turbulent** environment
- $L = 0.32 \text{ pc}$, $M_{\text{tot}} = 105 M_{\text{sol}}$
- Follow the collapse of the **densest** most massive region: $\sim 23 M_{\text{sol}}$
- Final resolution: $\sim R_{\text{sol}}$

Collapse with supersonic turbulence



- **Filament** with an attached sheet
- small **disk** within the filament (perpendicular)
- adiabatic (optically thick) core
- very efficient gas **accretion** through the filament

Mass accretion



- Very **high** mass accretion rates: up to $10 v_{\text{in}}^3/\text{G} \sim 10 M^3 c^3/\text{G}$ (cf. Shu '77 $dM/dt \sim c^3/\text{G}$)
- Mass accretion rates are higher than limits from radiation pressure by burning **massive** stars (e.g. *Wolfire & Cassinelli 1987*: $10^{-3} M_{\text{sol}}/\text{year}$)
- Protostars and disks assemble very **rapidly** within a supersonic turbulent environment

Ongoing & Future Projects

- Feedback from Jets \Rightarrow **Turbulence** from Jet-clump interactions?
- Clump/Core formation by colliding flows (*Vazques-Semandeni et al. 2006*)
- Include **Radiative Transfer** (Formation of massive stars; *Thomas Peters*)
- Precessing Jets (with *Christian Fendt*)

Summary

- Numerical simulations help to understand the star formation process
- Chemistry/**Cooling** important
- Magnetically driven **outflows** and **jets** from protostellar disk
- Magnetic fields reduce/prevent fragmentation
- **Turbulence** plays an important role (structure of cloud cores, accretion rates)
- No stage of hydrostatic cores in a turbulent environment: assembly of protostars is a very rapid process